

Spectral representation of lattice gluon- and ghost propagators

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Outline

1 Introduction

- Källén–Lehmann representation
- Tikhonov Regularization

2 Toy Models

- Breit-Wigner model
- $\int_0^\infty \rho(\omega)\omega \, d\omega = 0$ model

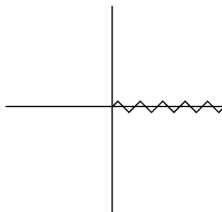
3 Lattice Data

- Gluon Propagator
- Ghost Propagator

4 Conclusion/Outlook

Introduction

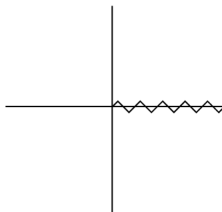
Introduction



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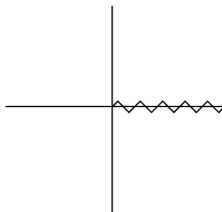


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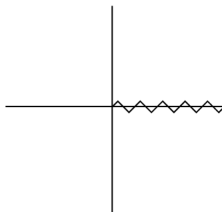


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- Analytically this has a unique analytical continuation to Minkowski $p_4^2 < 0$.
- Numerically this is very sensitive to noise: *ill-defined*.

Källén–Lehmann representation

Strategy: find the Källén–Lehmann representation.

$$G(p_4) = \int_{-\infty}^{\infty} \frac{\rho(\omega)}{\omega - ip_4} d\omega$$

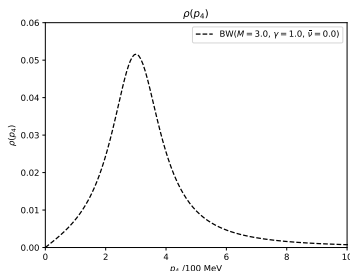
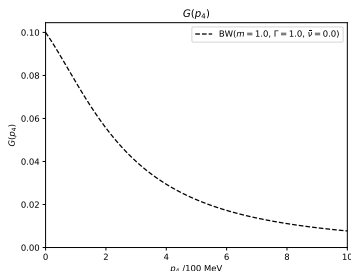
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Where $\rho(-\omega) = -\rho(\omega)$. The challenge: find $\rho(\omega)$ for given lattice data p_4 vs $G(p_4)$.



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- For observable quantities, $\rho(\omega) \geq 0$.
- Therefore, we expect positivity violations! Specifically:

$$\int_0^\infty \omega \rho(\omega) d\omega = 0$$

Inverting G : Naive approach

$$G(p_4) = \int_{-\infty}^{\infty} \frac{\rho(\omega)}{\omega - ip_4} d\omega = \int_{-\infty}^{\infty} K(p_4, \omega) \rho(\omega) d\omega, \quad K(p_4, \omega) := \frac{1}{\omega - ip_4}$$

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Problem

\mathbf{K} has small singular values, making it ill-conditioned on the numerical level \rightarrow ill-defined inversion problem

Sophisticated Approaches

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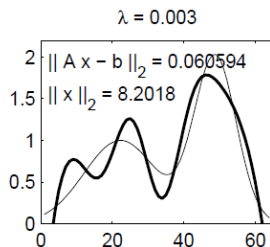
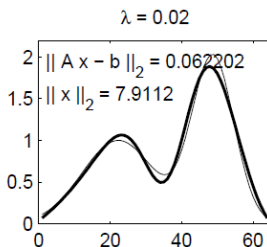
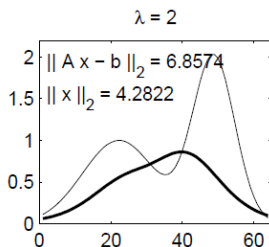
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- + Easy to implement
- No unique criterion to choose λ .

How to select the optimal λ ?



Too small lambda values mean no regularization at all: over-fitting. Too large results in over-smoothing.

Our preferred solution: Morozov's discrepancy principle

We have data G_n with standard deviation σ_n

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We therefore minimize

$$J_\lambda = \left\| \mathbf{K}\vec{\rho} - \vec{G} \right\|^2 + \lambda \|\vec{\rho}\|^2 \quad (2)$$

subject to $\left\| \mathbf{K}\vec{\rho} - \vec{G} \right\|^2 = \|\vec{\sigma}\|^2$. This has a unique solution.

Tikhonov Regularization

$$J_\lambda = \sum_{n=1}^N \frac{1}{\sigma_n^2} \left| \int_{-\infty}^{\infty} \frac{\rho(\omega)}{\omega - ip_n} d\omega - G_n \right|^2 + \lambda \int_{-\infty}^{\infty} \rho(\omega)^2 d\omega$$

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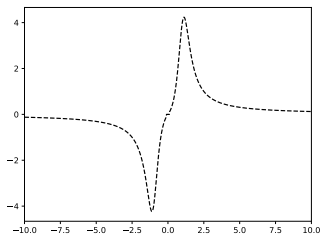
For a given λ

$$\rho(\nu) = -\frac{1}{\lambda} \sum_{n=1}^N \frac{1}{\sigma_n^2} \frac{c_n}{\nu - ip_n} \quad (3)$$

Toy Models

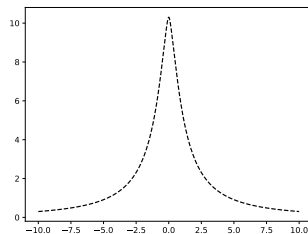
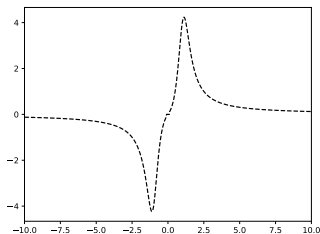
Reconstruction Process

$$\rho^{\text{orig}} \rightarrow G_n^{\text{orig}} \rightarrow G_n \in \mathcal{N}(G_n^{\text{orig}}, (\epsilon G_n^{\text{orig}})^2)$$



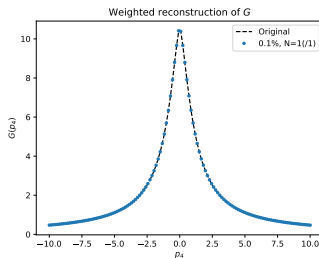
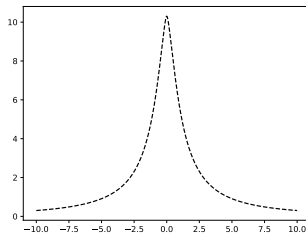
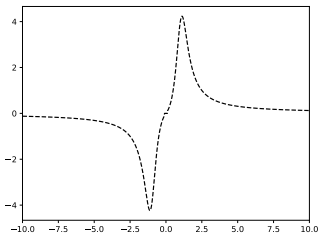
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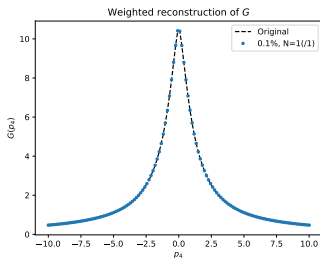
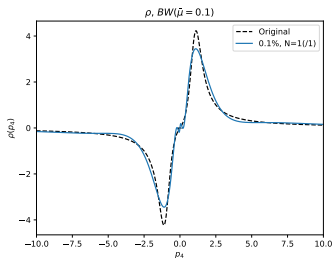
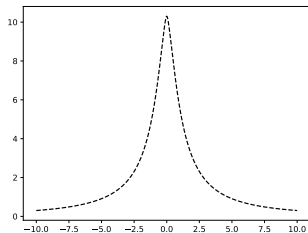
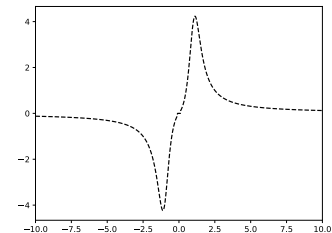
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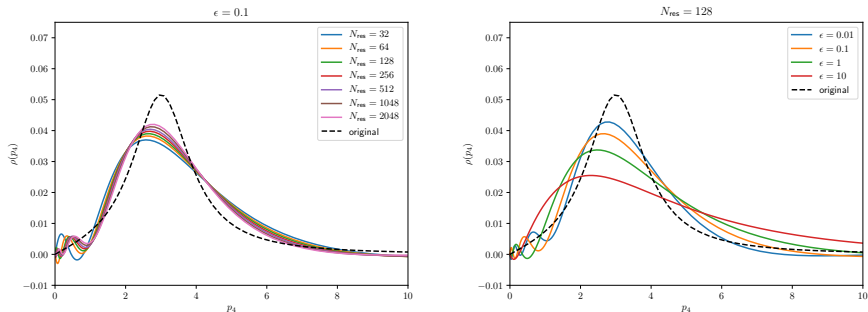


Figure: $BW(M = 300\text{MeV}, \gamma = 100\text{MeV}, \bar{\nu} = 0)$

$$\rho(\omega) = \frac{1}{\pi} \frac{2\omega\gamma}{(\omega^2 - \gamma^2 - M^2)^2 + 4\omega^2\gamma^2}$$

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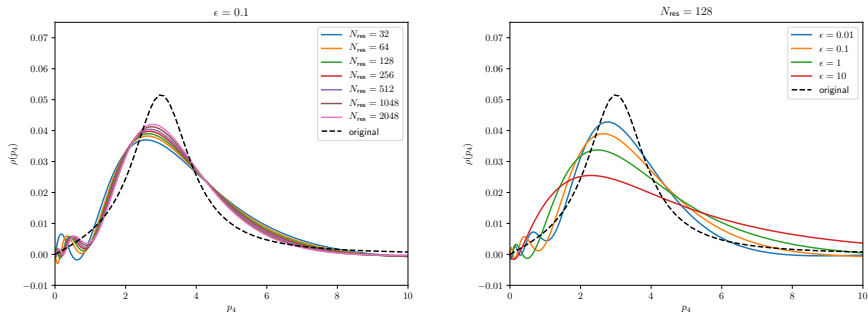


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Integrals are preserved within statistical error with Tikhonov reconstruction

Quality of the fit

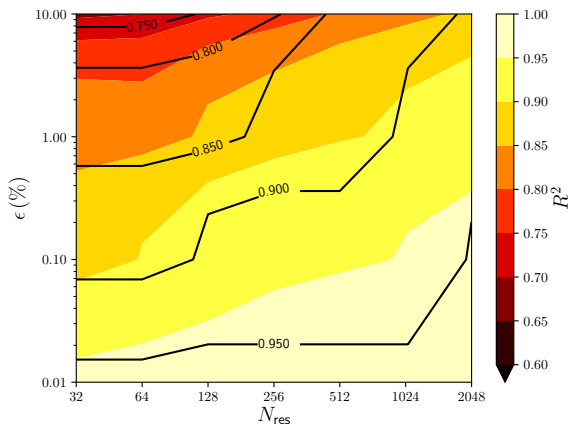


Figure: R^2 for the quality of the fit

$$\int_0^\infty \rho(\omega) \omega \, d\omega = 0 \text{ model}$$

For gluons we expect the following sum rule:

$$\int_0^\infty \rho(\omega) \omega \, d\omega = 0 \quad (4)$$

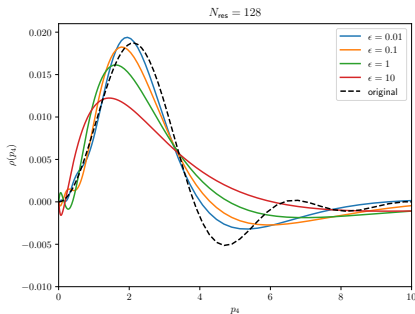
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Lattice Data

Observation

$T = 0$ Gluon propagator data, 80^4 lattice, $\beta = 6.0$, normalised at 3GeV

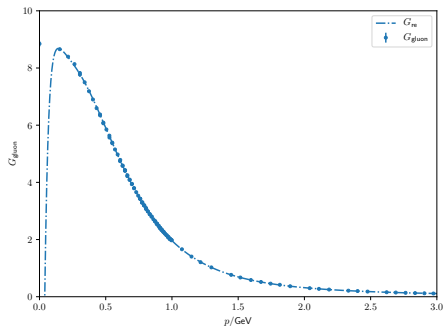


Figure: Gluon Propagator

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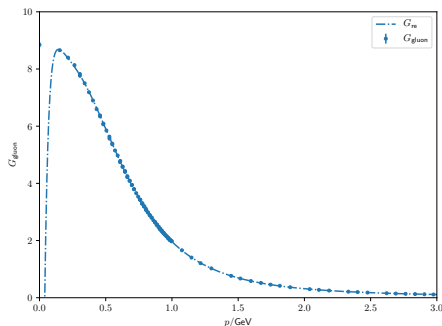


Figure: Gluon Propagator

- The datapoint at $p_4 = 0$ is very poorly reconstructed.
- Pragmatic solution:

$$\int_{-\infty}^{\infty} \rightarrow \left(\int_{-\infty}^{-\omega_0} + \int_{\omega_0}^{\infty} \right)$$

ω_0 VS α

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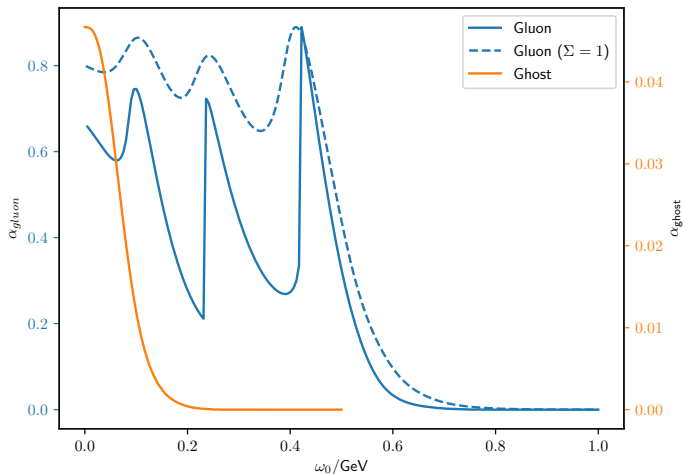


Figure: ω_0 vs α

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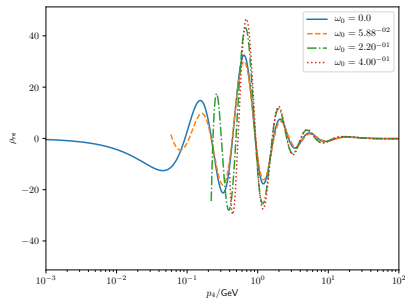
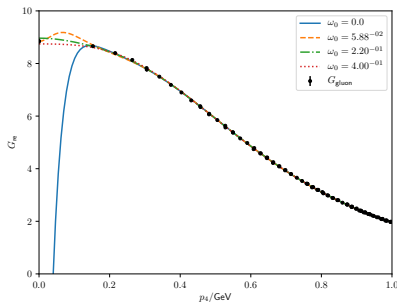


Figure: Left: Gluon propagator. Right: Gluon spectral density

Ghost Propagator

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- Numerical inversions don't like δ -peaks.
- It can be avoided by inverting $g(p) = p^2 G(p)$ instead.

Ghost Propagator

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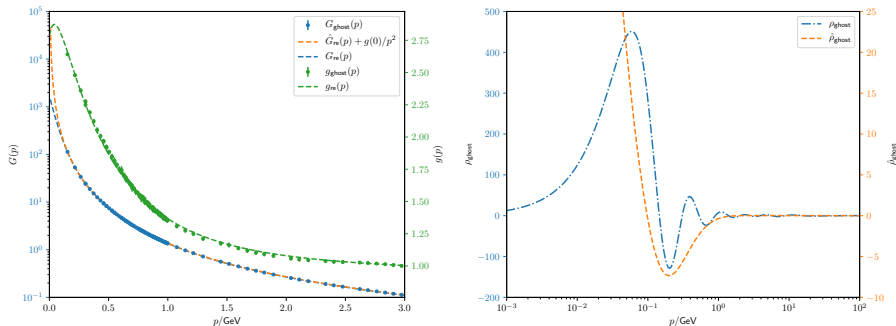


Figure: Left: ghost propagator. Right: ghost spectral representation

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- + Gluons and ghosts both display positivity violations, as they should
- + Study $T > 0$ datasets, especially around the critical temperature

Dudal, David, Orlando Oliveira, and Paulo J. Silva (2014).

“Källén-Lehmann spectroscopy for (un)physical degrees of freedom”. In: *Phys. Rev. D* 89.1, p. 014010. DOI: 10.1103/PhysRevD.89.014010. arXiv: 1310.4069 [hep-lat].

Tripolt, Ralf-Arno et al. (2018). “Numerical analytic continuation of Euclidean data”. In: arXiv: 1801.10348 [hep-ph].

Roelfs et al. 2018?

ip -Formalism, weighted vs unweighted

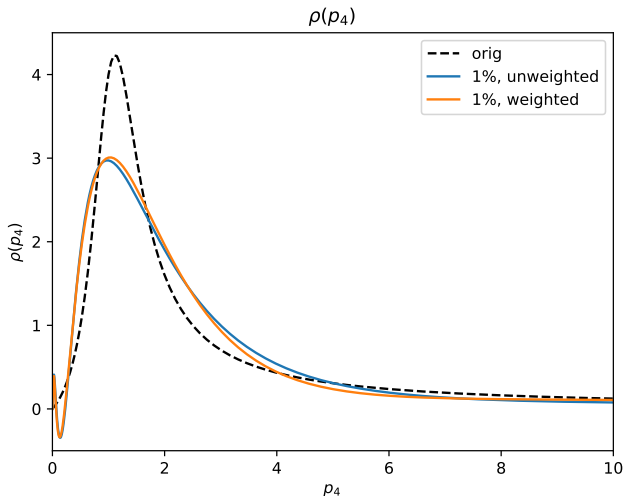


Figure: Weighted vs unweighted, resp. $N_p = 256$

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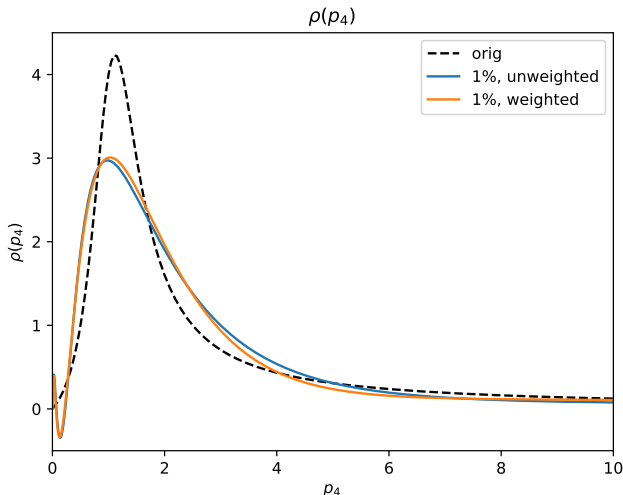


Figure: Weighted vs unweighted, resp. $N_p = 256$

Weighted performs better in the UV, use the weighted algorithm from now

Reproducibility/Stability

Does every sample $G_n \in \mathcal{N}(G_n, (\epsilon G_n)^2)$ give approximately the same λ ?

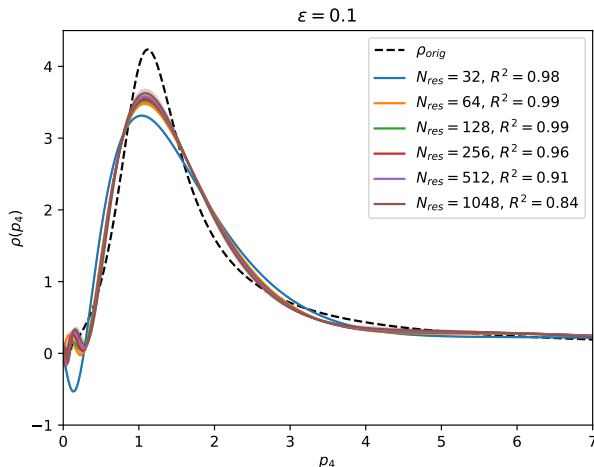


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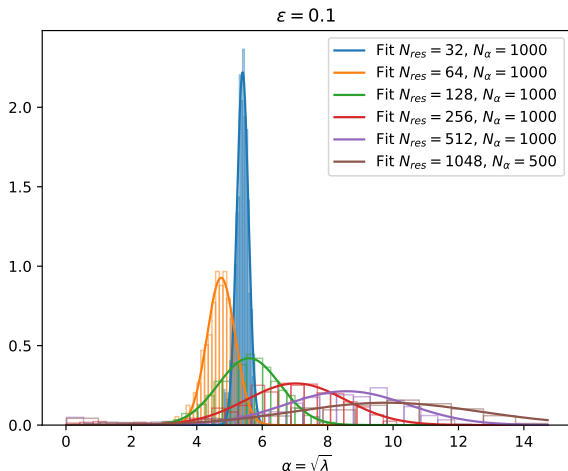


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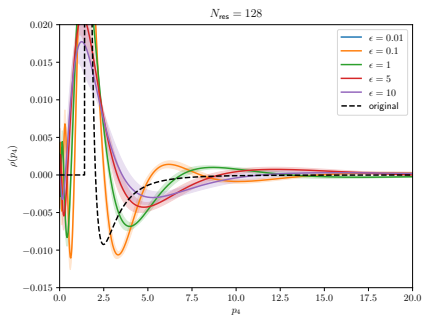
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Toy model

$$\rho(\omega) = \begin{cases} \frac{-1}{\omega^4+4} + \frac{A}{\omega^6+2} \\ \rho(-\omega) = -\rho(\omega) \end{cases} \quad \omega > \sqrt{2}$$



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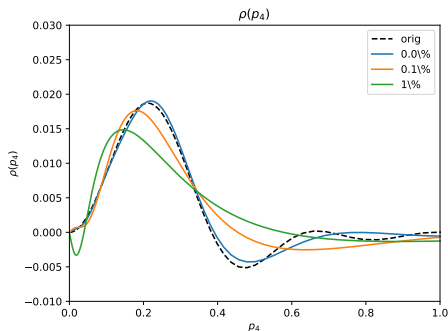
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$\int \rho(\omega) \omega d\omega = 0$ model, ip vs p^2

